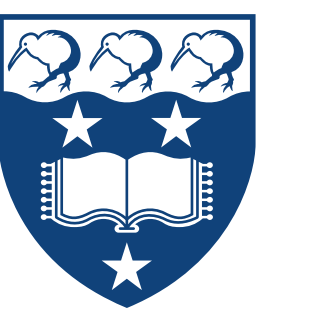




# PHASE RESETTING IN THE YAMADA MODEL OF A Q-SWITCHED LASER

J. Ngaha<sup>1,2</sup>, N. G. R. Broderick<sup>2,3</sup>, and B. Krauskopf<sup>1,2</sup>

<sup>1</sup>Department of Mathematics, The University of Auckland, New Zealand  
<sup>2</sup>The Dodd-Walls Centre for Photonic and Quantum Technologies, Dunedin, New Zealand  
<sup>3</sup>Department of Physics, The University of Auckland, New Zealand



## MOTIVATION – A WAY TO PROBE STABLE OSCILLATING SYSTEMS

- Optical frequency combs, optical clocks, and optical neural networks all work with stable oscillating optical systems.
- After an induced perturbation, the system will return back to stable oscillations.
- With phase resetting, we can study how the new phase of the return relates to the phase of the initial oscillation.

## THE YAMADA MODEL

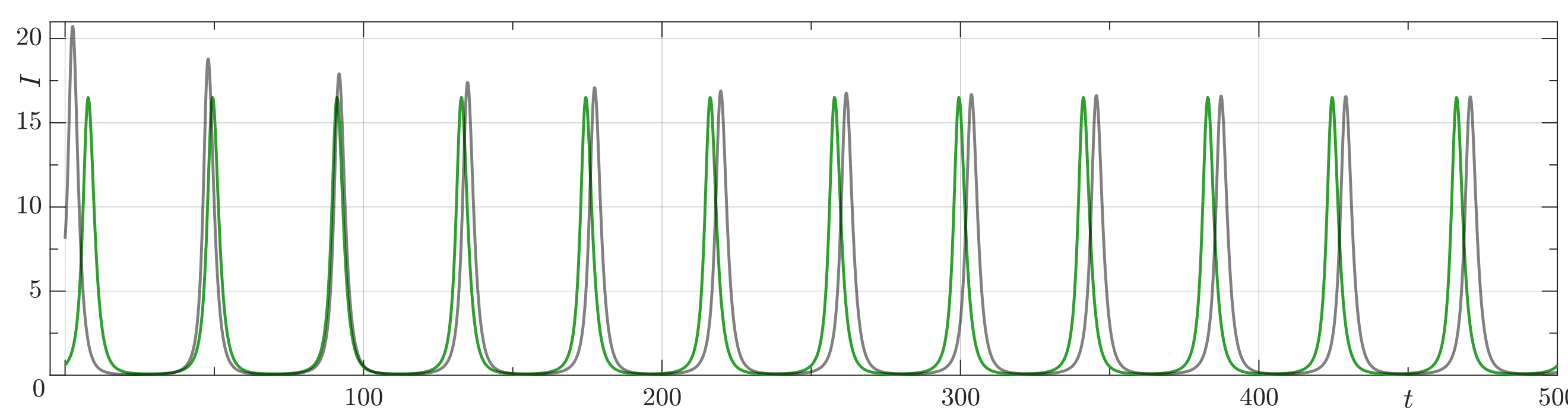
- Yamada model for a q-switched laser with a lossy saturable absorber:

$$\begin{aligned}\dot{G} &= \gamma(A - G - GI) \\ \dot{Q} &= \gamma(B - Q - aQI) \\ \dot{I} &= I(G - Q - 1)\end{aligned}$$

where:

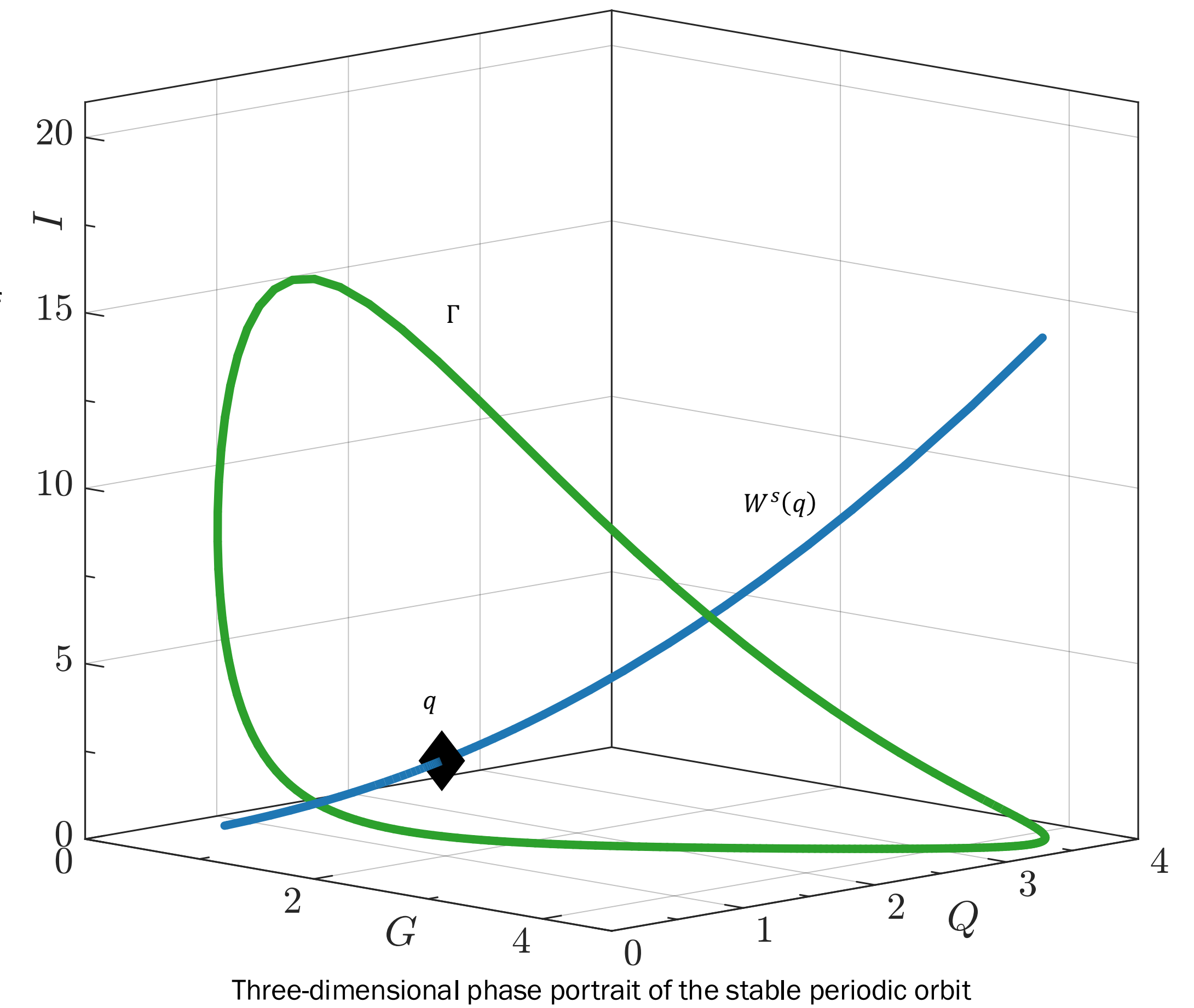
- $G$  – gain,
- $Q$  – absorption,
- $I$  – laser intensity,
- $\gamma = 0.0354$  – photon loss rate,
- $A = 7.3757$  – pump current,
- $B = 5.8, a = 1.8$  – fixed, dimensionless parameters.

- Nine different regions of dynamics [1,2], separated by different bifurcations. Here we focus on a region with a stable periodic solution.



Time series of the unperturbed oscillations (green) and perturbed response (black) back to the original pulsations at  $t \sim 300$ , with a phase shift.

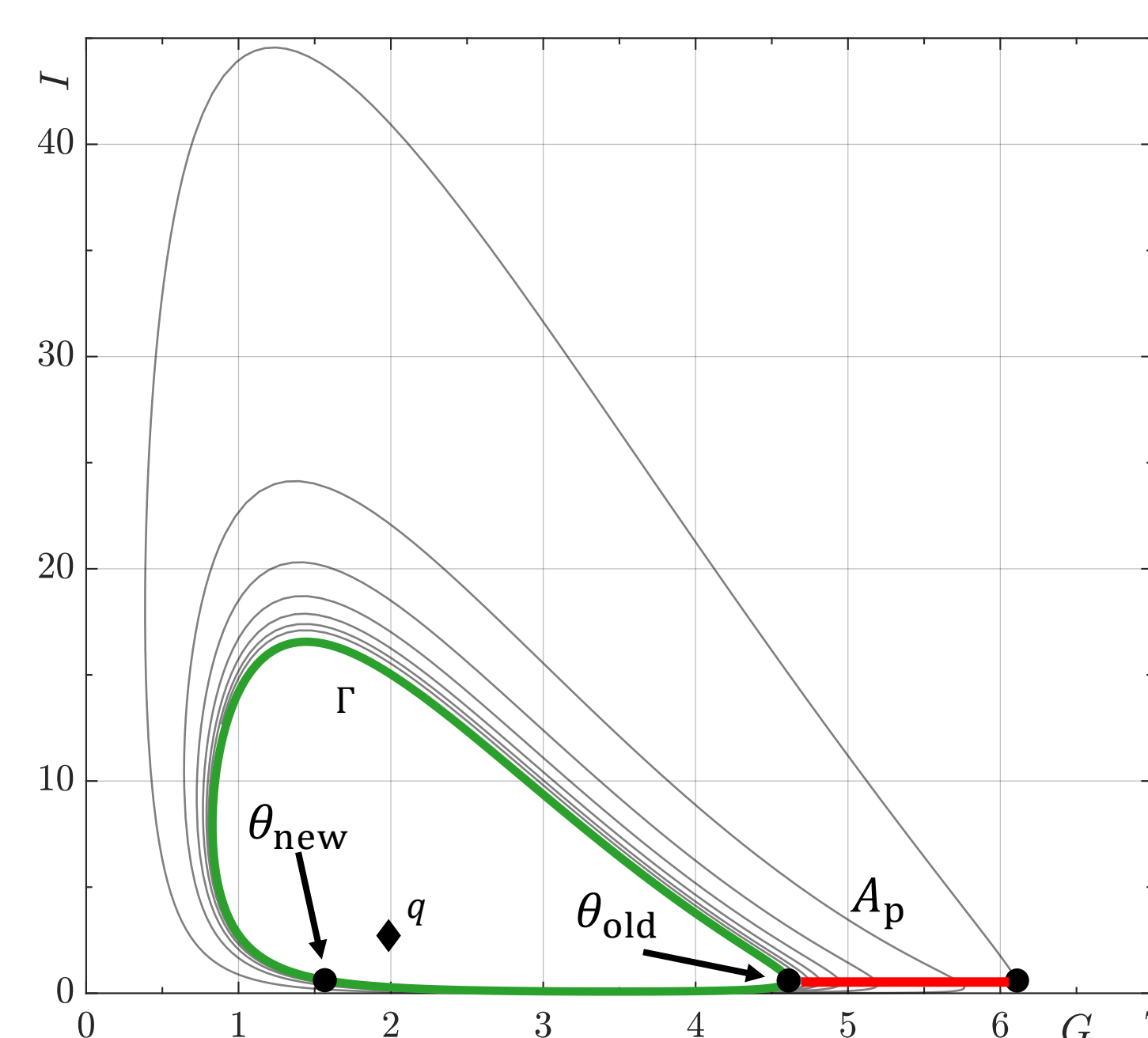
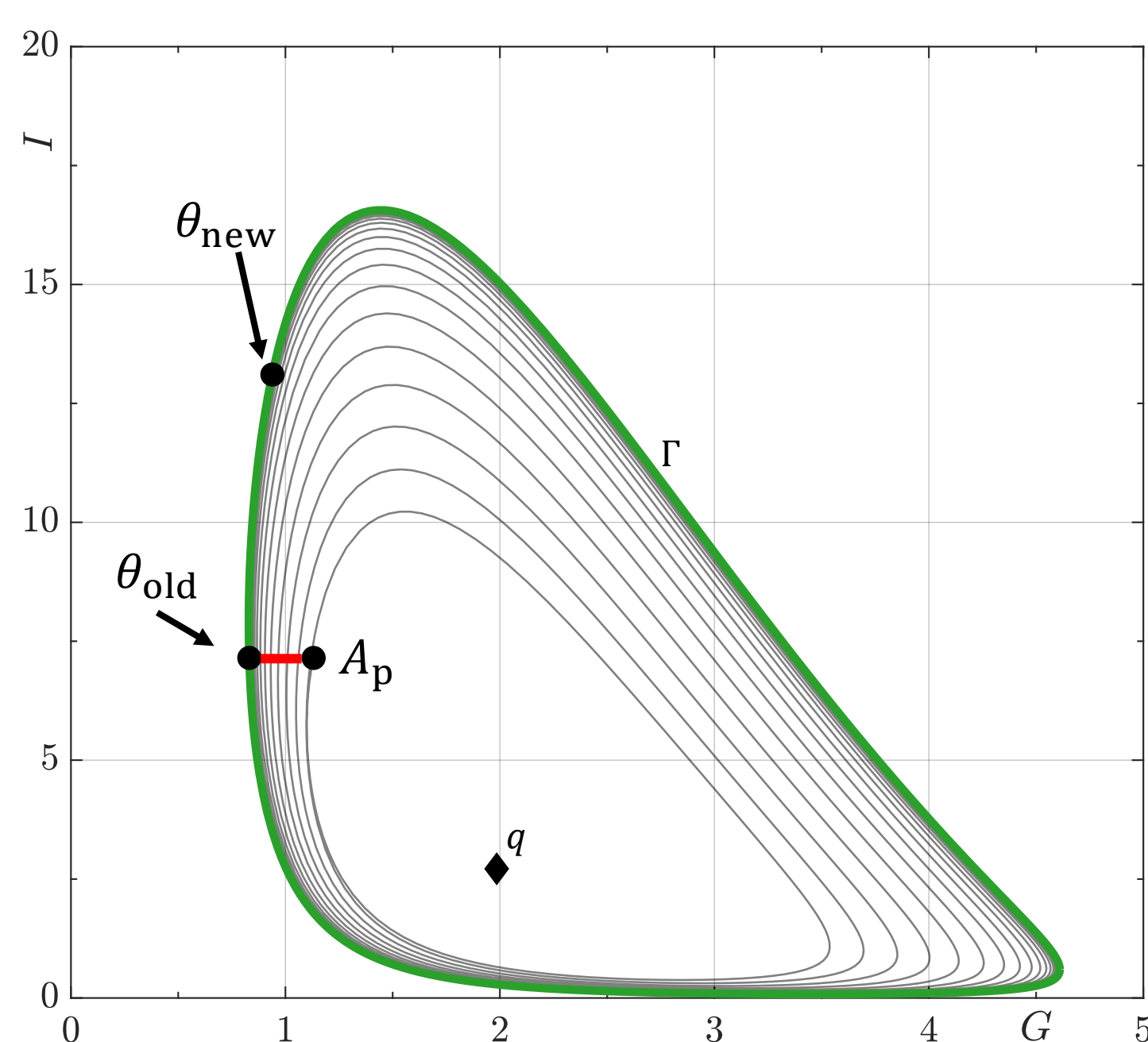
- Objects in phase space:
  - $o$  – stable equilibrium (not shown),
  - $q$  – unstable stationary point (diamond),
  - $\Gamma$  – stable periodic orbit (green curve),
  - $W^s(q)$  – one-dimensional stable manifold of  $q$  (blue curve).
- All initial conditions not on  $W^s(q)$  will evolve towards the periodic orbit  $\Gamma$ .



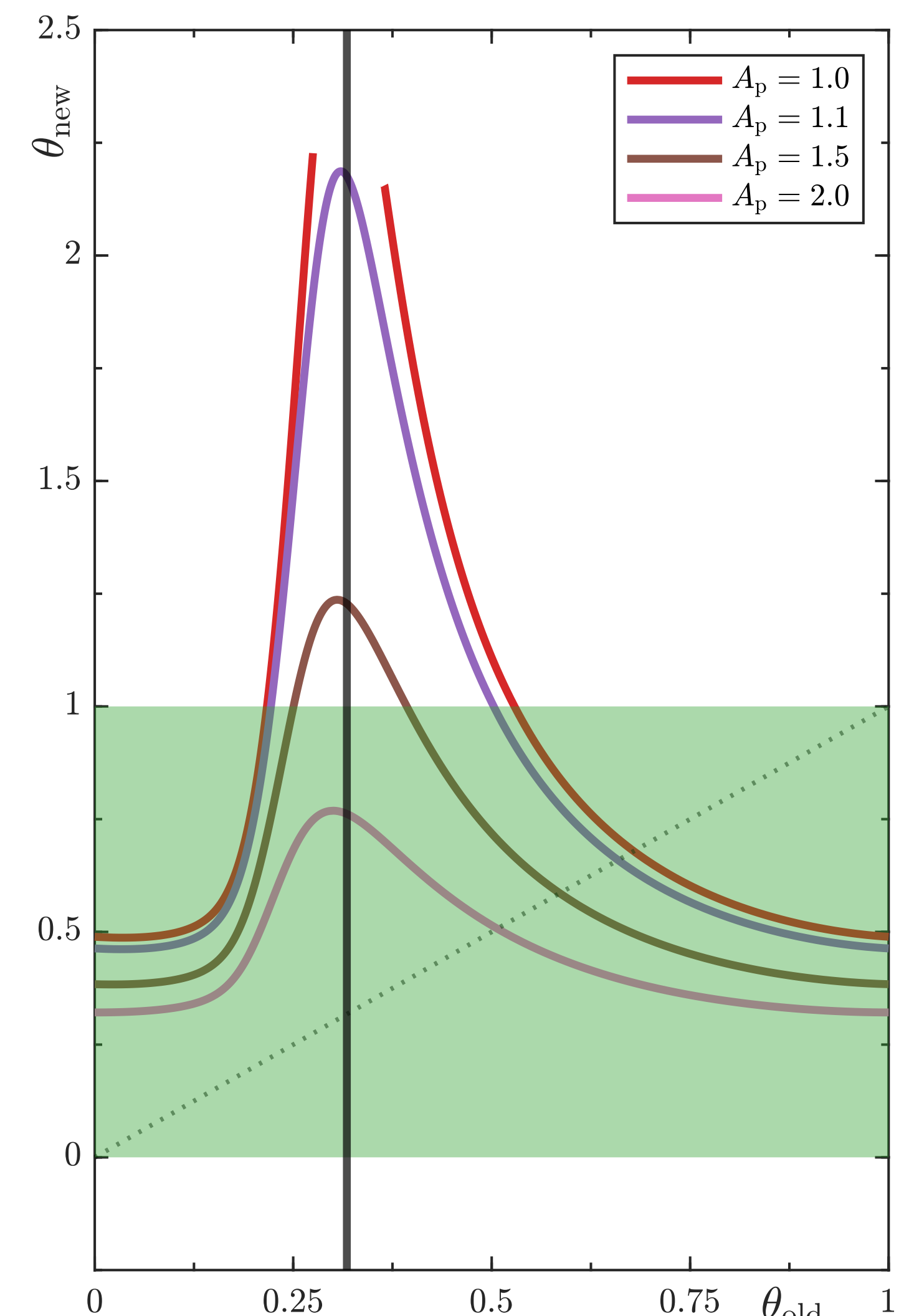
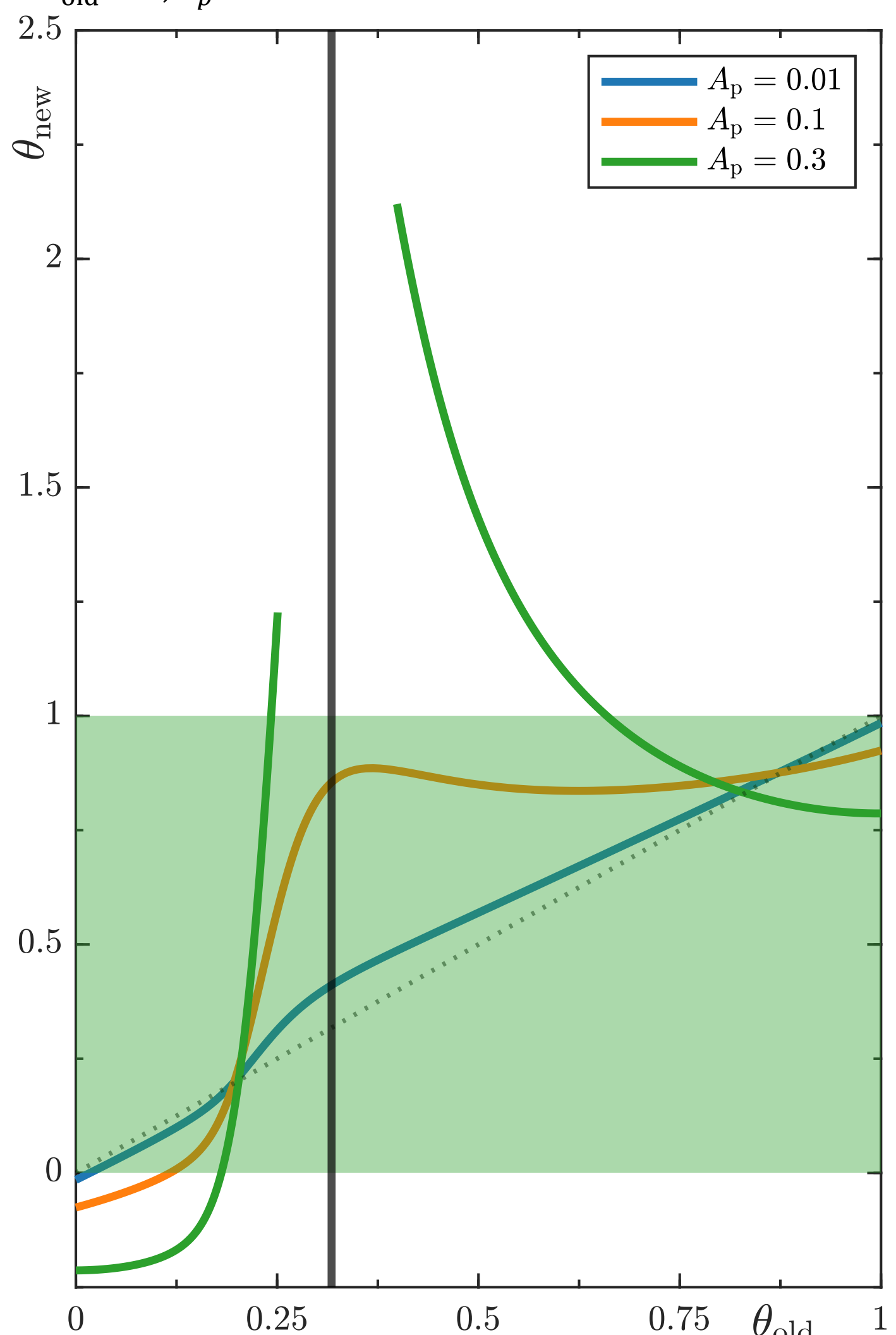
Three-dimensional phase portrait of the stable periodic orbit

## PHASE RESETTING AND PHASE TRANSITION CURVES (PTCs)

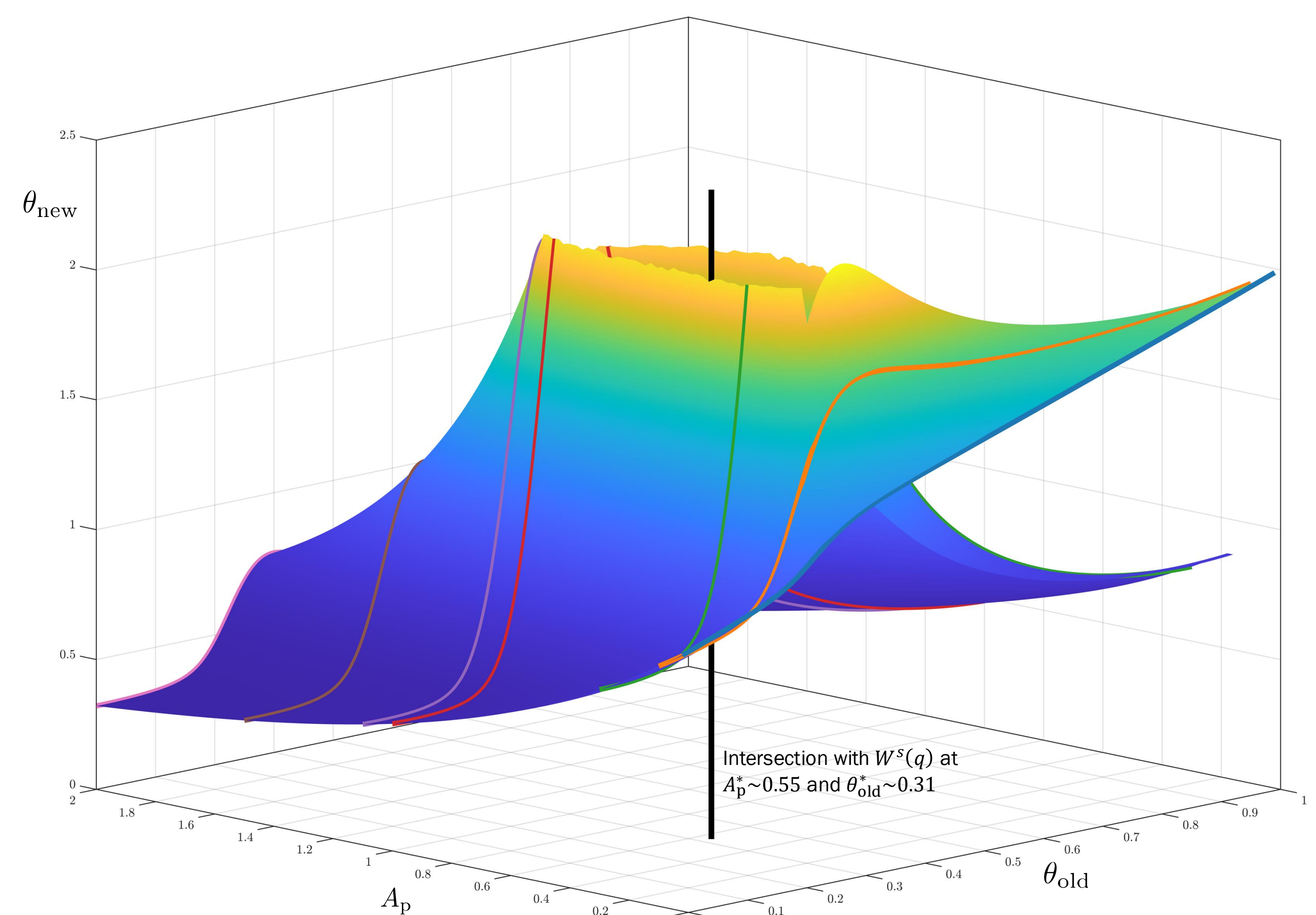
- Perturbations cause a phase shift ('lag') in intensity pulses [3,4].
- Phase resetting parameters for fixed perturbations in positive  $G$ -direction:
  - $A_p$  – perturbation amplitude,
  - $\theta_{old}$  – phase where perturbation is applied,
  - $\theta_{new}$  – phase where trajectory “resets”.
- For weak perturbations,  $\theta_{new}$  increases as  $\theta_{old}$  increases.
- For strong perturbations, there is a topological change.
- For perturbation  $A_p^* \sim 0.55$  at  $\theta_{old}^* \sim 0.31$ ,  $\Gamma$  intersects  $W^s(q)$ . The trajectory does not return to  $\Gamma$ , and the phase is undefined.



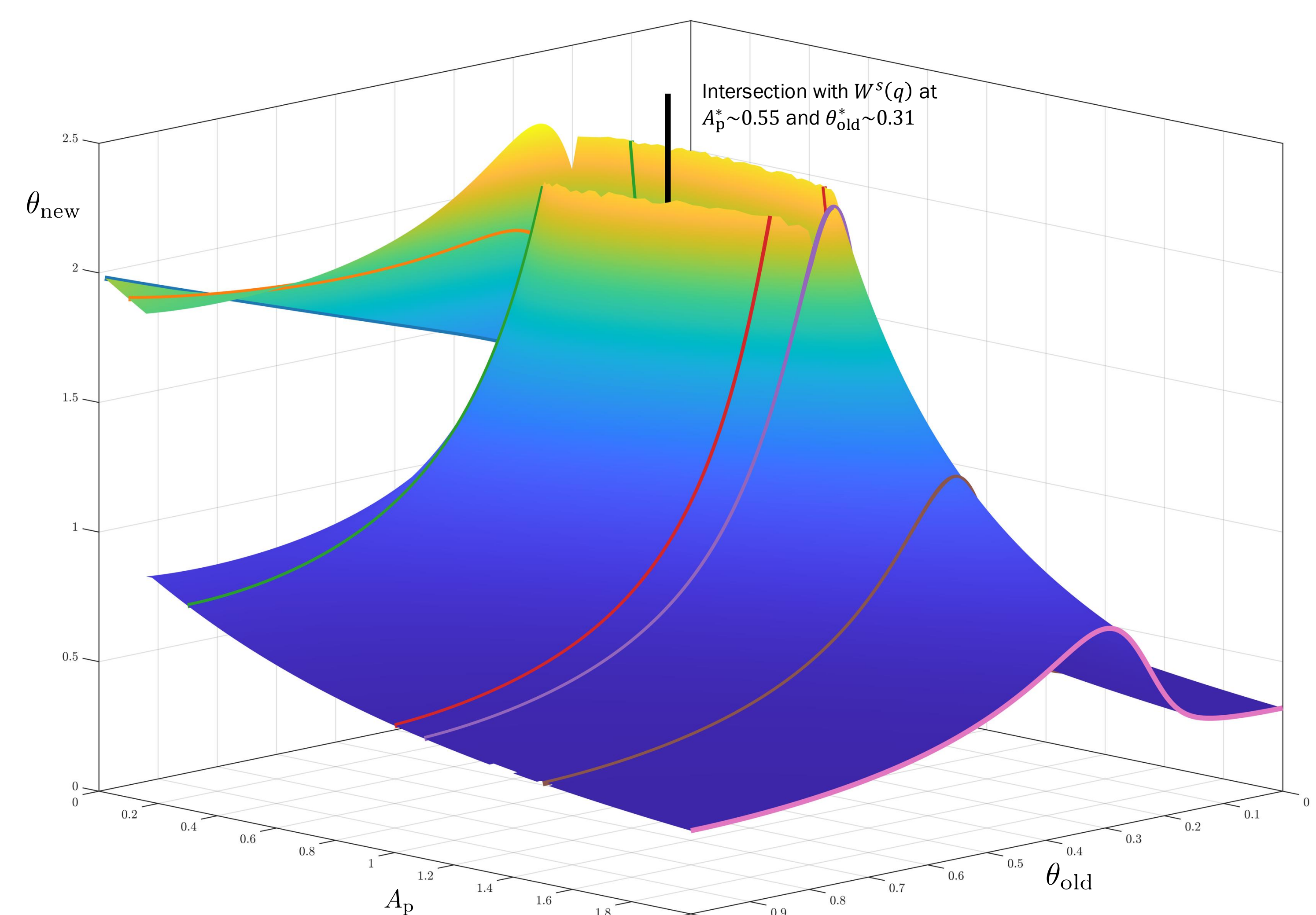
Phase portraits of the periodic orbit with a perturbation in the positive  $G$ -direction at: (left)  $\theta_{old} = 0.5$  with  $A_p = 0.3$ ; and at (right)  $\theta_{old} = 0, A_p = 1.5$ .



PTCs for different perturbation amplitudes in the positive  $G$ -direction: (left) before the intersection with  $W^s(q)$ , and (right) after the intersection with  $W^s(q)$ . The phase at which  $\Gamma$  intersects  $W^s(q)$  is  $\theta_{old}^* \sim 0.31$  (black line).



PTC surface for perturbations in the positive  $G$ -direction.



## References

- M. Yamada, “A theoretical analysis of self-sustained pulsation phenomena in narrow-stripe semiconductor lasers”. *IEEE J. Quantum Electron* **29**, 1330 (1993).
- J. L. A. Dubbeldam and B. Krauskopf. “Self-pulsations of lasers with saturable absorbers: Dynamics and bifurcations”. *Opt. Commun.*, **159**, 325 (1999).
- P. Langfield, B. Krauskopf, and H. M. Osinga. “A continuation approach to computing phase resetting curves” in *Advances in Dynamics, Optimization, and Computation*, Vol 304 of SSSC (Springer International Publishing, 2020), pp. 3-30.
- K. H. Lee, N. G. R. Broderick, B. Krauskopf, and H. M. Osinga. “Phase response to arbitrary perturbations: Geometric insights and resetting surfaces”. *Discrete and Continuous Dynamical Systems-B* (2024).